

Quantum Mechanics in Quantum Reference Frame

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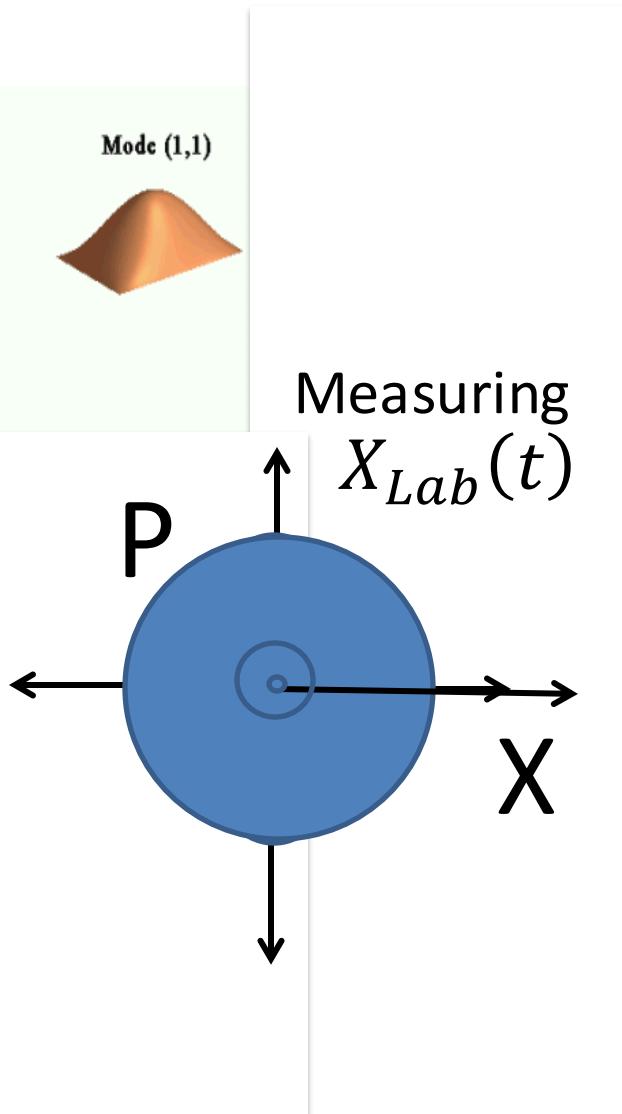


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Quantum limits of measurement of motion and fields

$$X_{Lab}(t) = X \sin(\omega t) + P \cos(\omega t)$$



$$\text{Var}(X) \text{Var}(P) \geq 1/4 \quad [X, P] = i$$

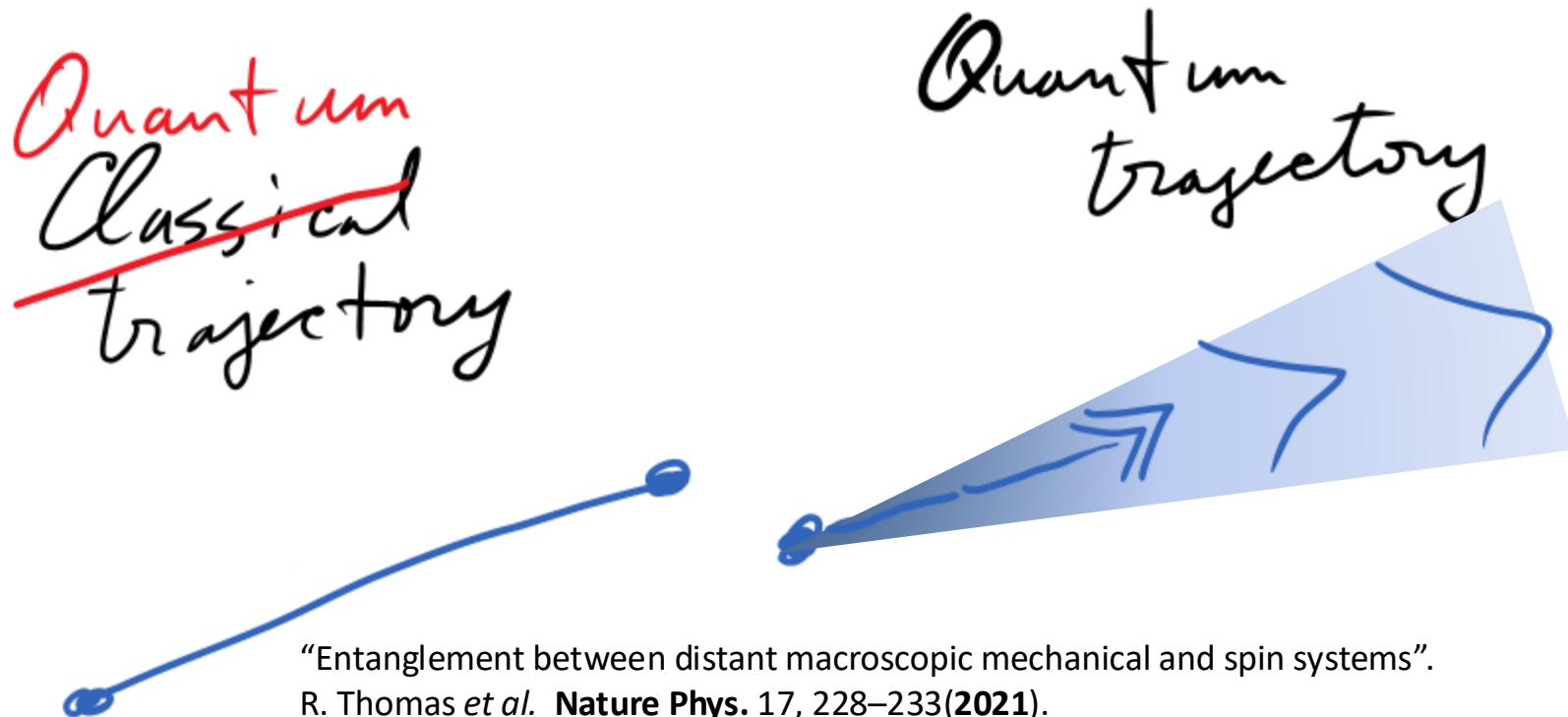
Measurement of X imposes uncertainty on P
and the other way around:

Quantum Back Action of measurement

And yet, arbitrary small perturbations in BOTH position and momentum can be measured simultaneously

Trajectories without quantum uncertainties in a negative mass reference frame

AKA *quantum-mechanics-free subspaces*



"Entanglement between distant macroscopic mechanical and spin systems".

R. Thomas *et al.* **Nature Phys.** 17, 228–233(2021).

"Quantum back-action-evading measurement of motion in a negative mass reference frame"

C. Moller *et al.*, **Nature** 22980 (2017)

"Establishing Einstein-Podolsky-Rosen channels between nanomechanics and atomic ensembles". K. Hammerer, M. Aspelmeyer, ESP, P. Zoller. **PRL** 102, 020501 (2009).

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See also Tsang and Caves. Quantum mechanics free subspaces, **PRL** 2010.

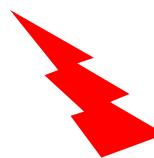
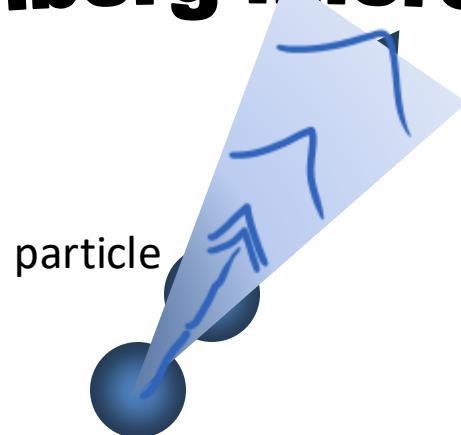


Standard quantum limit example: displacement measurement

“Heisenberg microscope”



W. Heisenberg



photon

$$X(t) = X + \frac{Pt}{m}, \quad \Delta X \Delta P \geq \frac{\hbar}{2} \Rightarrow$$
$$[\Delta X(t)]^2 \geq (\Delta X)^2 + \frac{\hbar^2 t^2}{4m^2(\Delta X)^2} \geq \frac{\hbar t}{m} \quad (\text{SQL})$$

Measurement of motion beyond SQL in a negative mass reference frame

1. Define trajectory relative to a quantum reference
2. Reference system has an effective negative mass
3. Entangled state of the reference and the probed systems is generated
4. Measurement (relative to reference frame)

“Entanglement between distant macroscopic mechanical and spin systems”.

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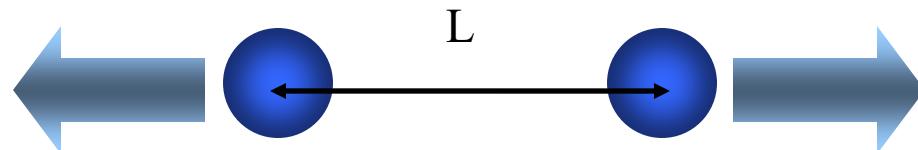
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Einstein-Podolsky-Rosen (EPR) entanglement 1935

2 particles entangled
in position/momentum

$$[\hat{X}_1 - \hat{X}_2, \hat{P}_1 + \hat{P}_2] = 0$$



$$\hat{X}_1, \hat{P}_1 \quad \hat{X}_2, \hat{P}_2$$

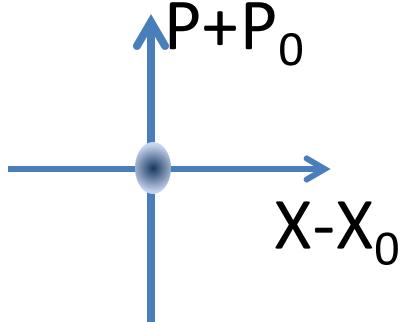
$$\hat{X}_1 - \hat{X}_2 = L \quad \hat{P}_1 + \hat{P}_2 = 0$$

Simon (2000); Duan, Giedke, Cirac, Zoller (2000)

Necessary and sufficient condition for entanglement

$$Var(X - X_0) + Var(P + P_0) < 2$$

Trajectory in a quantum reference frame with negative mass



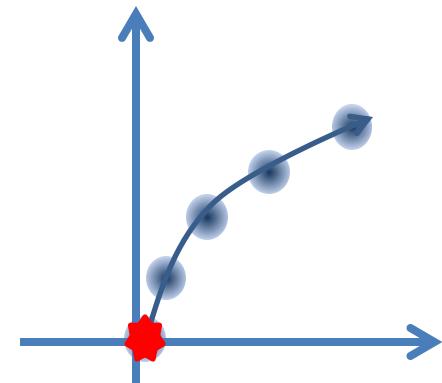
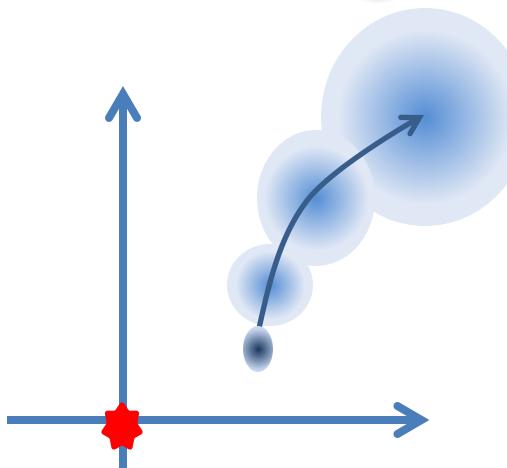
$$X - X_0 = X_{X_0} \rightarrow 0$$
$$P + P_0 \rightarrow 0$$

Probe system entangled with origin system

$$X(dt)_{X_0} = X(0)_{X_0} + (\dot{X} - \dot{X}_0)dt$$
$$= X(0)_{X_0} + (P + P_0)dt$$

$$m = -m_0 = 1$$

Not good enough



$$X(t) = X(0) \cos(\omega t) + P(0) \sin(\omega t)$$

Joint measurement of motion of two oscillators, one with negative mass and frequency ($\omega = -\omega_R$):

Reference oscillator with negative mass, frequency, and k

$$\omega^2 = k/m$$

$$X(t) - X_R(t) = [X(0) - X_R(0)] \cos(\omega t) + [P(0) \cancel{+} P_R(0)] \sin(\omega t)$$

EPR:

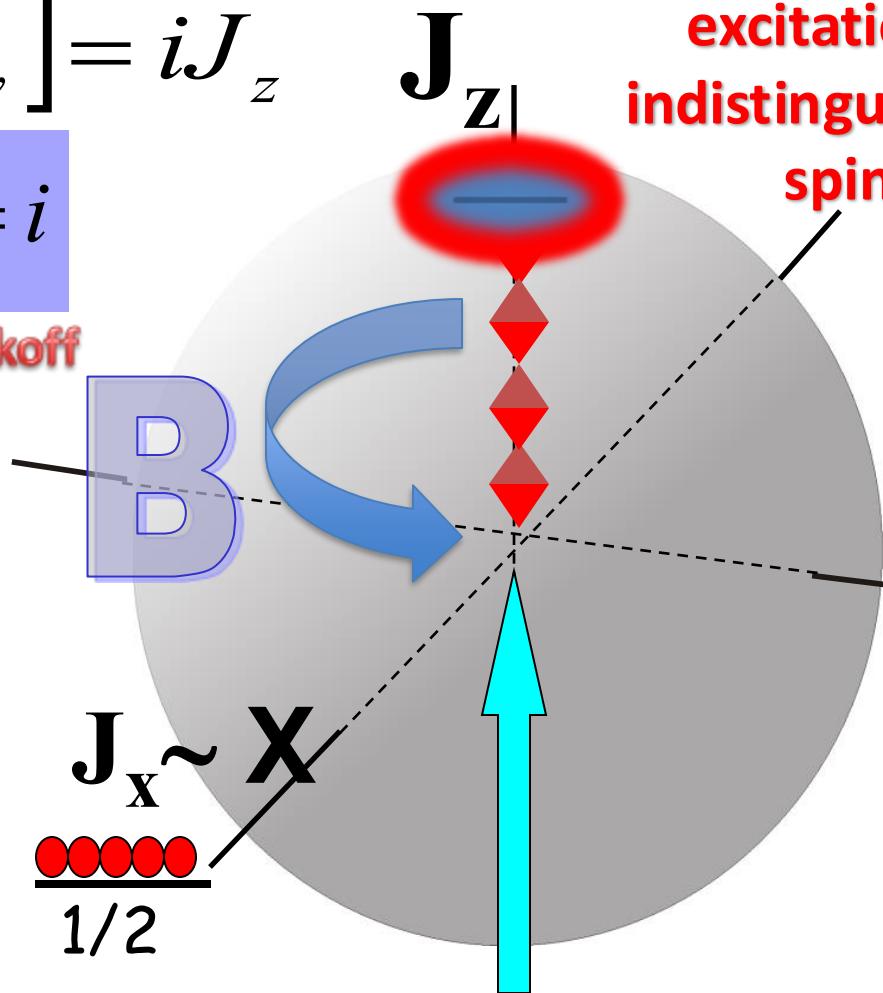
$$Var(X - X_0) + Var(P + P_0) < 2$$

Spin ensemble = positive/negative mass oscillator

$$[\hat{J}_x, \hat{J}_y] = i\hat{J}_z$$

$$[\hat{X}, \hat{P}] = i$$

Holstein-Primakoff



-1/2

$$\left| -\frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \dots \right\rangle + \left| \frac{1}{2}, -\frac{1}{2}, \frac{1}{2}, \dots \right\rangle + \left| \frac{1}{2}, \frac{1}{2}, -\frac{1}{2}, \dots \right\rangle + \dots$$

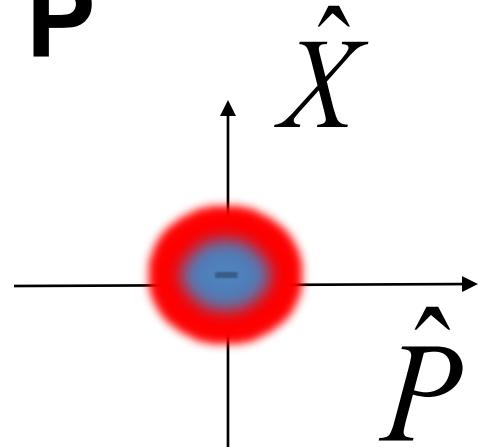
Collective
excitation of
indistinguishable
spins

$$J = \sum_{i=1}^N j_i$$

Oscillator with negative
mass, frequency, and k

$$\omega^2 = k/m$$

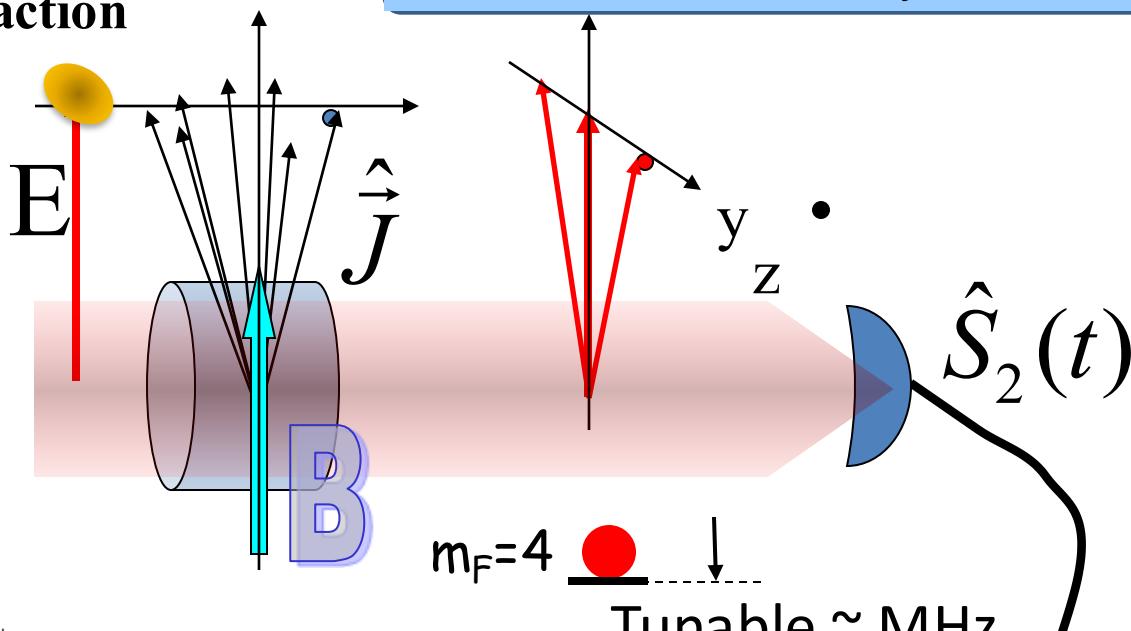
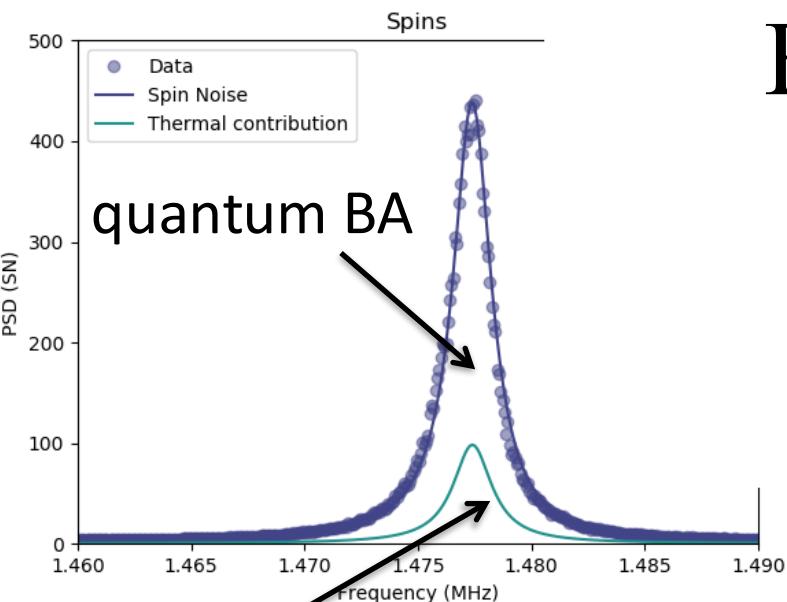
$$\mathbf{J}_y \sim \mathbf{P}$$



Julsgaard, Kozhekin, ESP, Nature 2001
Dideriksen et al. Nature Comm. 2021

Quantum noise of measurement of spin oscillator: probe noise and backaction

$$J_z^{lab} = J_z^{rot} \cos \Omega t - J_y^{rot} \sin \Omega t$$



(Almost)
Ground
Spin state

$$X_{\text{spin}} \sim J_z^{\text{rot}}(t) + \text{Back action noise}$$

$$P_{\text{spin}} \sim J_y^{\text{rot}}(t) + \text{Back action noise}$$

Cooperativity

Susceptibility

$$P_{L1,out} = -P_{L1,in} + \text{spin state} + \Gamma_S \chi_S X_{L1,in}$$

Shot noise of light

Back action of light

Application of measurement in negative mass reference frame #1

Quantum limited and entanglement assisted
magnetometry with 10^{-16} Tesla/ $\sqrt{\text{Hz}}$ sensitivity

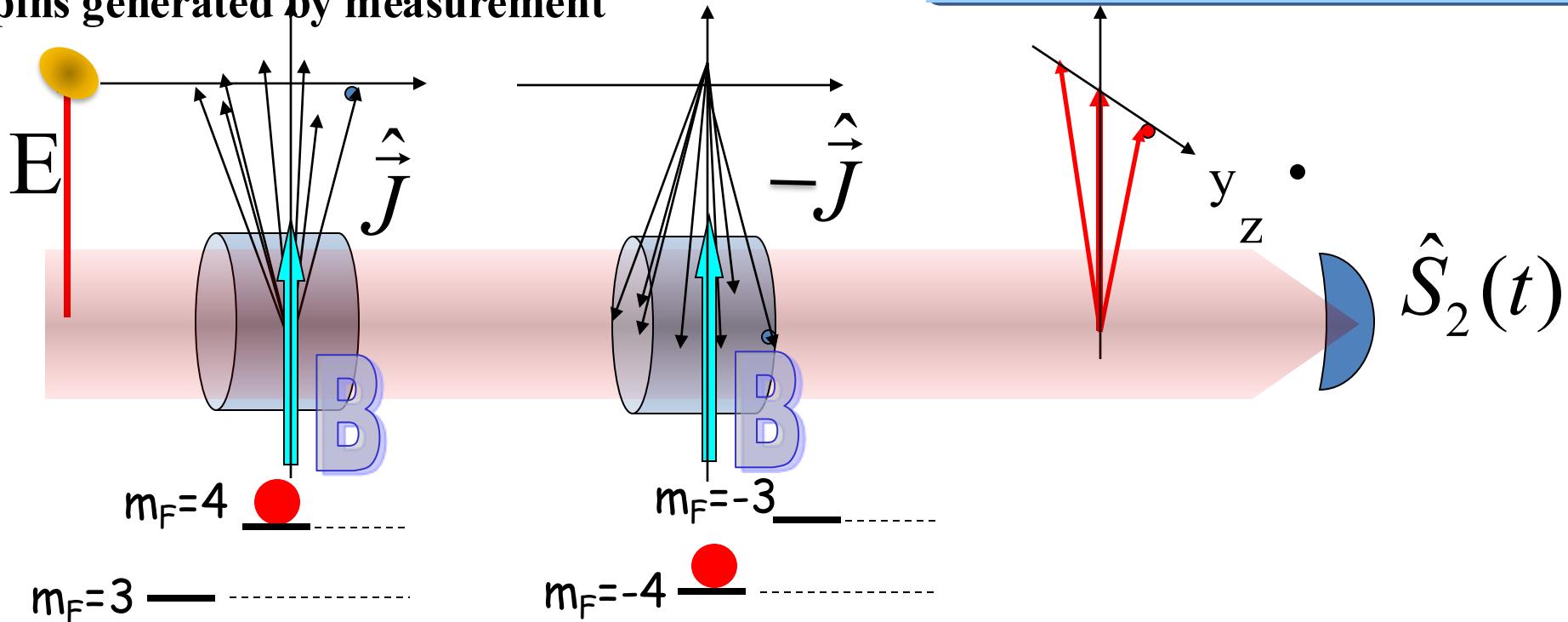
- Two magnetic sensors with opposite spin orientations
- Generation of entanglement of two sensors
- Magnetic field applied and measured

Phys.Rev.Lett. 104, 133601 (2010) [arXiv:0907.2453](https://arxiv.org/abs/0907.2453)

[W. Wasilewski](#), [K. Jensen](#), [H. Krauter](#), [J.J. Renema](#), [M. V. Balabas](#), E.P.

Quantum entanglement of two macroscopic spins generated by measurement

$$J_z^{lab} = J_z^{rot} \cos \Omega t - J_y^{rot} \sin \Omega t$$



$$P_{L,out} = -P_{L,in} + \sqrt{\Gamma_S \gamma_S} \chi_S (J_{z1}^{rot} + J_{z2}^{rot}) + \Gamma_S \chi_S X_{L1,in}$$

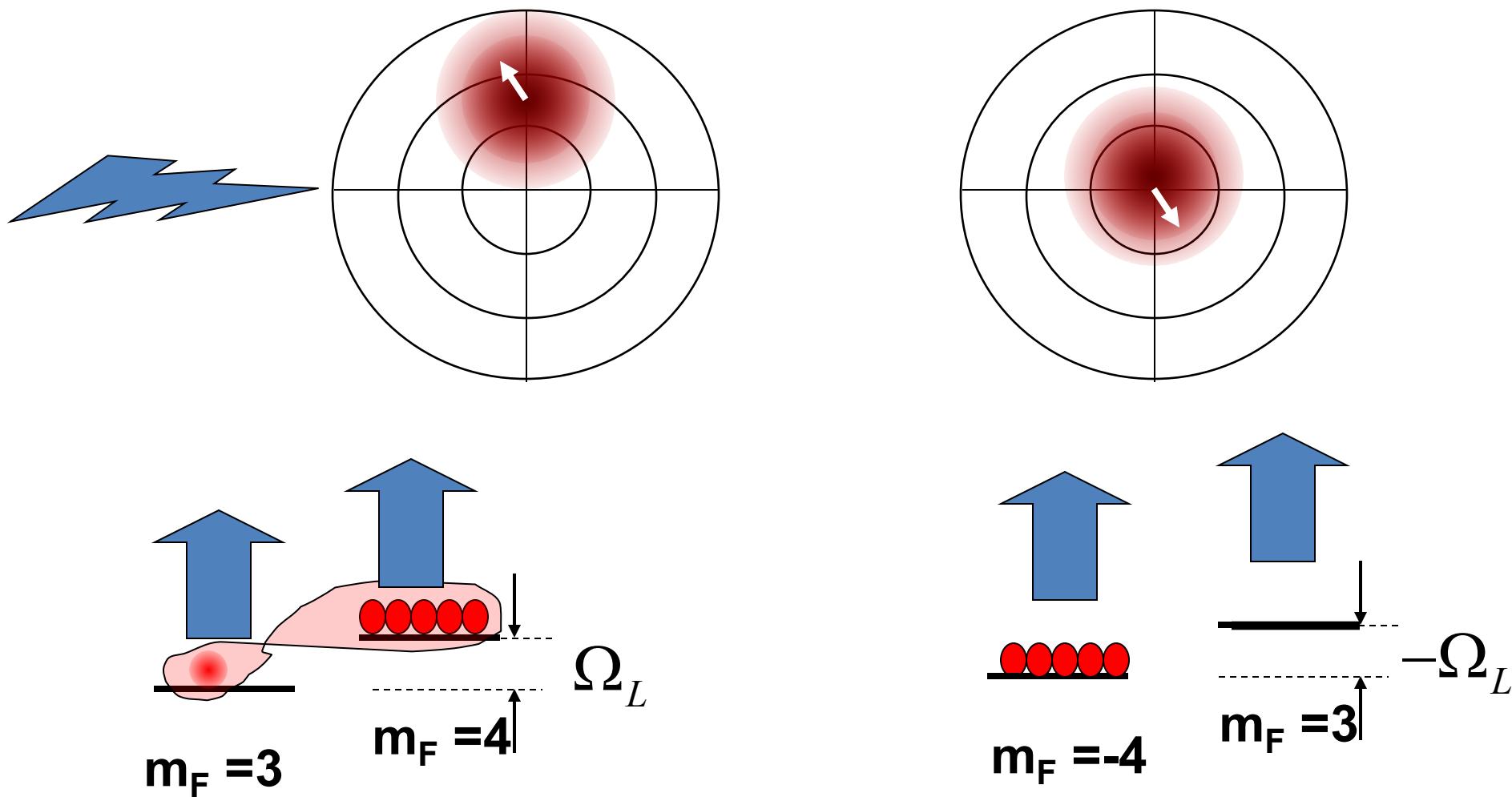
Shot noise of light

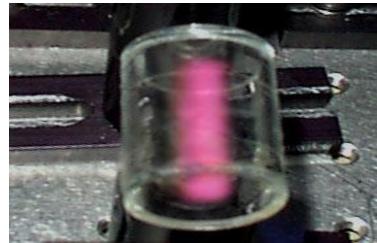
Back action of light cancelled out in the measurement

Equal cooperativities

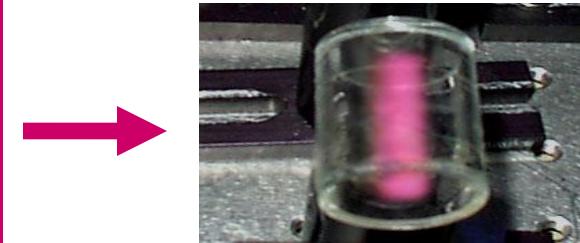
Opposite sign susceptibilities

Quantum back action of probe light on atoms: cancellation via entanglement of two ensembles



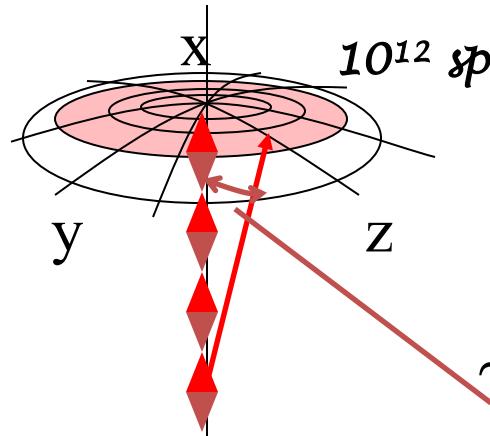


*Entanglement
of two
macroscopic
objects.*



$$Var(\hat{J}_{z1} + \hat{J}_{z2})/2J_x + Var(\hat{J}_{y1} + \hat{J}_{y2})/2J_x < 1$$

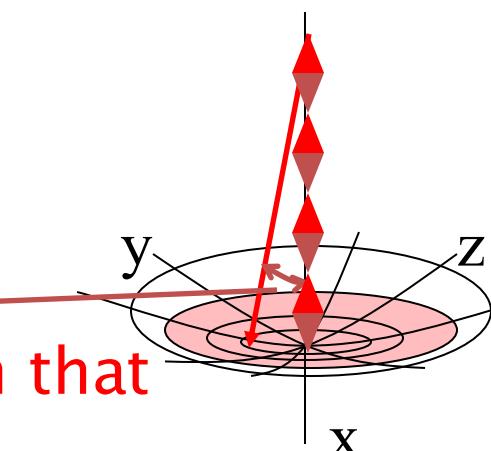
Can be created by a measurement



10¹² spins in each ensemble

$$\sim N^{-\frac{1}{2}}$$

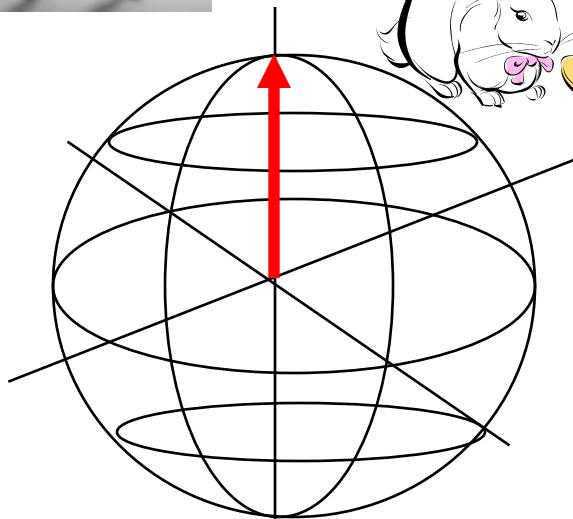
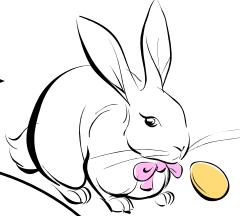
Spins which are “more parallel” than that are entangled



Detection of tiny oscillating magnetic fields



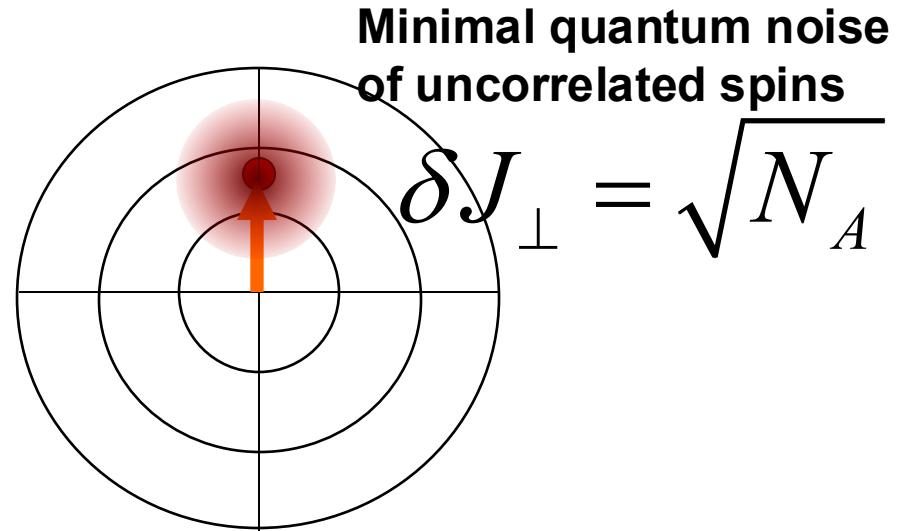
$$\vec{B}_{RF} = b \cos(\Omega t)$$



$$\vec{B}$$

Bias magnetic field
Larmor frequency Ω

Spin dynamics
top view



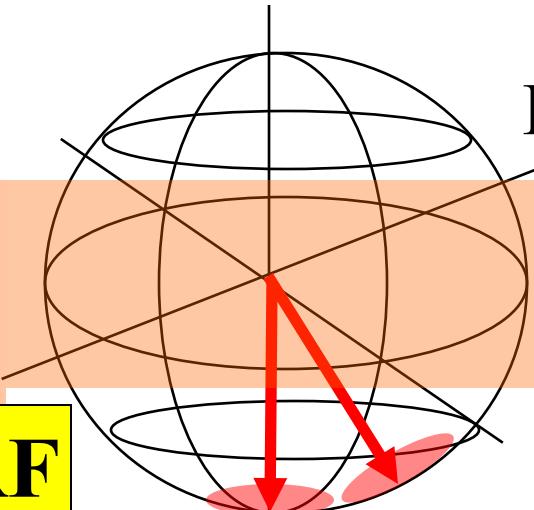
$$J_{\perp} \approx \gamma B_{RF} N_A T_2$$

γ – Gyromagnetic constant
 T_2 – Transverse spin coherence time

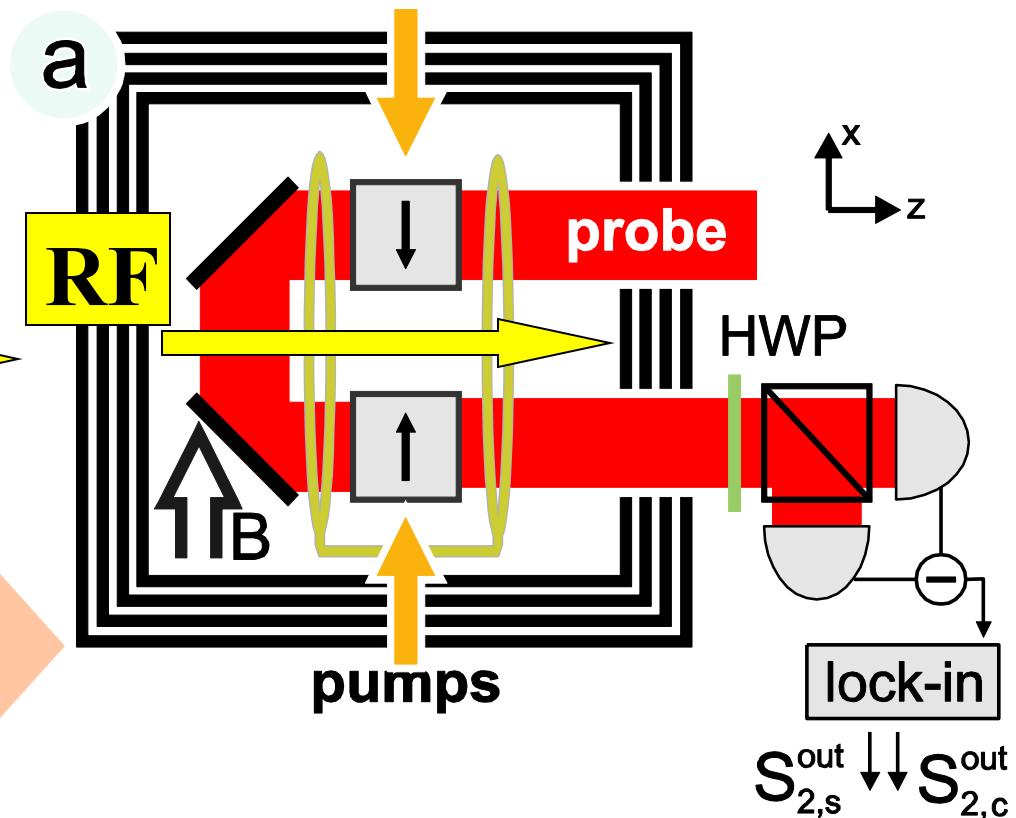
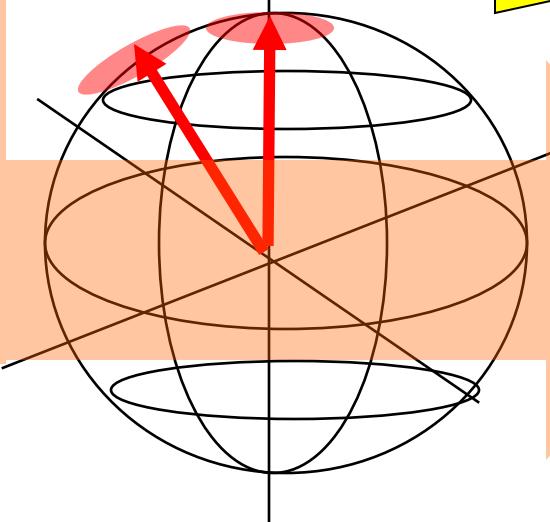
Calcellation of measurement back action with two cells

$$\left[\hat{J}_{z1} + \hat{J}_{z2}, \hat{J}_{y1} + \hat{J}_{y2} \right] = i(\hat{J}_{x1} + \hat{J}_{x2}) = 0$$

measurement does not change relative spin orientation

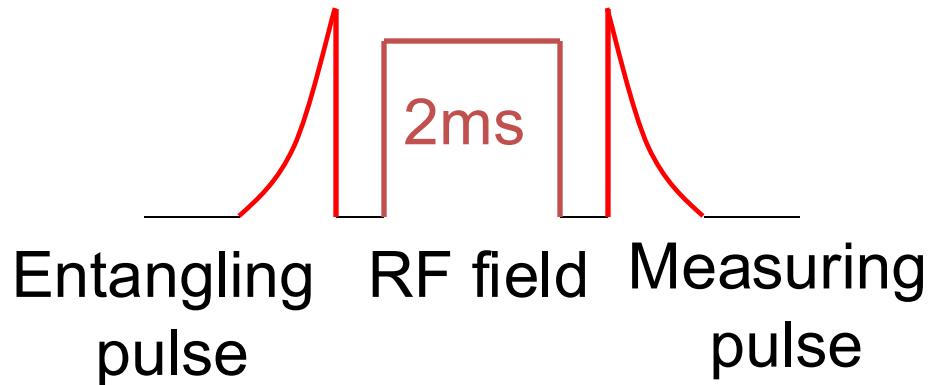


Faraday rotation signal doubles



Magnetometry beyond the projection noise limit

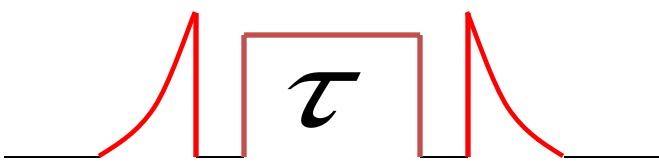
Entanglement by QND measurement



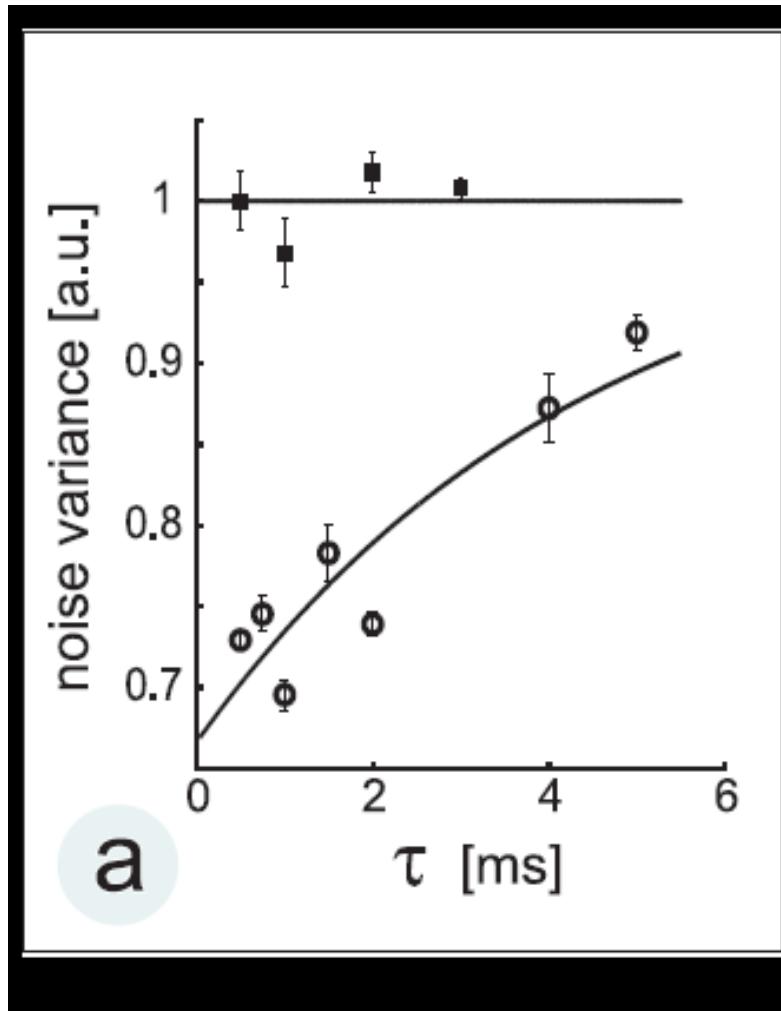
$$\delta\left(\hat{J}_{z1} + \hat{J}_{z2}\right)^2 + \delta\left(\hat{J}_{y1} + \hat{J}_{y2}\right)^2 = 1.3 \cdot J_x < 2 \cdot J_x$$

$$\delta\left(\hat{X}_1 + \hat{X}_2\right)^2 + \delta\left(\hat{P}_1 + \hat{P}_2\right)^2 = 1.3 < 2$$

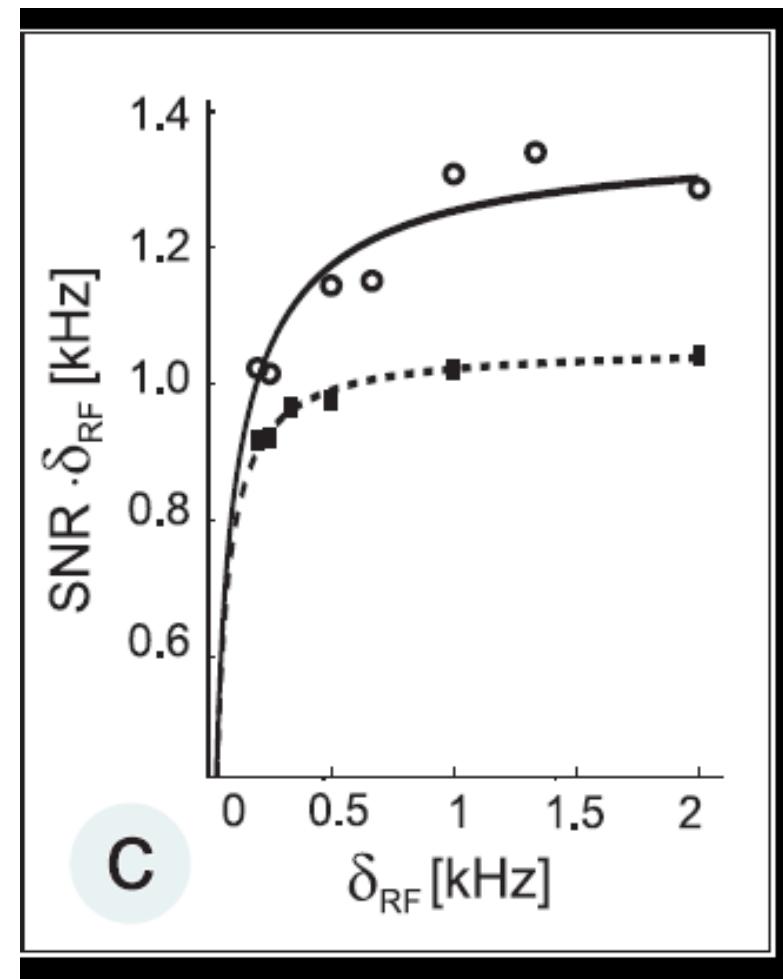
EPR entanglement



Signal/noise times
bandwidth

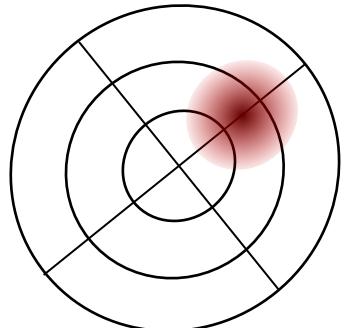


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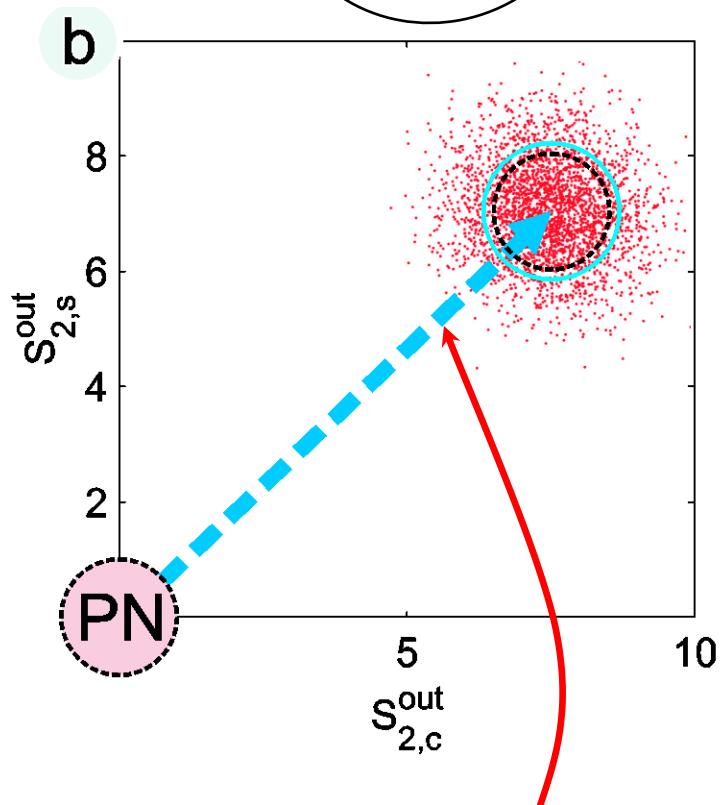


c

Magnetic field sensitivity with $1.5 \cdot 10^{12}$ atoms



$$0.42 \cdot 10^{-15} T / \sqrt{Hz}$$



State-of-the-art cell magnetometer
with 10^{16} K atoms

Lee et al, Appl. Phys. Lett. 2006

$$0.24 \cdot 10^{-15} T / \sqrt{Hz}$$

100-fold improvement
in sensitivity per atom

$$B_{RF} = 36 \cdot 10^{-15} \text{ Tesla} = 3.6 \cdot 10^{-10} G$$